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# A Reference-Feed-Forward-Based Damping Method for Virtual Synchronous Generator Control

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**Abstract**—The virtual synchronous generator (VSG) control has been extensively applied for realizing grid-friendly interconnections of power converters. Nevertheless, due to the direct shaft dynamics emulation, low-frequency power oscillations may be introduced in the case of limited damping effects. To solve this issue, power control law modifications have been proposed based on the feedback control theory. Although power oscillations can be suppressed to an acceptable level, the original inertial response of VSG will be significantly degraded. This paper proposes a new damping method based on the reference feed forward (RFF) control. The proposed method is able to effectively attenuate poorly-damped power oscillations in VSG-controlled converters without affecting the original inertial response. Theoretical analysis and effectiveness of the proposed damping method have been validated by both EMT simulations and experiments.

**Index Terms**—Virtual synchronous generator (VSG), power oscillation, oscillation damping, reference feed forward control

## I. INTRODUCTION

VSG has been applied in various applications to enable islanded operation, inertial response and frequency support [1]–[3]. However, the virtual inertia embedded in VSGs may introduce low-frequency active power oscillations, which subsequently affects the system stability and increase the risk of damaging power electronic converters [4], [5].

In order to suppress the low-frequency active power oscillations, various damping algorithms have been developed [6]–[10]. In [6], additional damping effects have been realized by using the frequency slip estimated from a phase-locked loop (PLL). One problem of using this PLL-based method is the PLL-related stability issue, especially when grids are weak [11]. Additionally, without the help of PLLs, a damping correction control approach was proposed in [7], where the algorithm was developed based on modifying the active power feedback. Although acceptable attenuation of power oscillations can be attained by adjusting the filter passband, the degrees of freedom are limited to achieve fully adjustable power dynamics. Benefiting from the flexibility of state-feedback control, the method proposed in [8] enables fully controllable

power dynamics. Moreover, the damping approach derived from the relationship between power dynamics and angular accelerations was proposed in [9]. It is able to provide fully controllable power dynamics for individual VSG as well as parallel connected VSGs. In [10], another approach was proposed by adding a virtual damping element.

One common disadvantage of aforementioned methods is that the damping term based on the feedback control inevitably affects the original inertial response of VSG [12]. Specifically, when loads change rapidly, the rate of change of frequency (RoCoF) can not be limited by inertia as desired [12]. To solve this issue, this paper proposes a new damping method using the feed forward control, and its technical contributions can be summarized as: i) No degradation of the inertial response; ii) Fully controllable damping effects and power dynamics; iii) Enable automatic tuning and real-time adaptation, in the case that system parameters are measured or estimated.

The rest of the paper is organized as follows. In Section II, the small-signal model is presented. In Section III, power oscillations and existing damping algorithms are analyzed. Section IV introduces the proposed RFF-based damping method. In Section V, simulation and experimental results are given, which leads to the final conclusion in Section VI.

## II. MODELING OF THE VSG-CONTROLLED SYSTEM

Fig. 1 indicates the general scheme of a VSG, where  $R_f$  and  $L_f$  are converter-side resistor and inductor, and  $C_f$  denotes the filter capacitor. The impedance of the interconnection transformer is denoted by  $Z_t$ , the grid impedance is represented by  $Z_g$ , and  $Z_L = Z_t + Z_g$  is the total grid-side impedance.

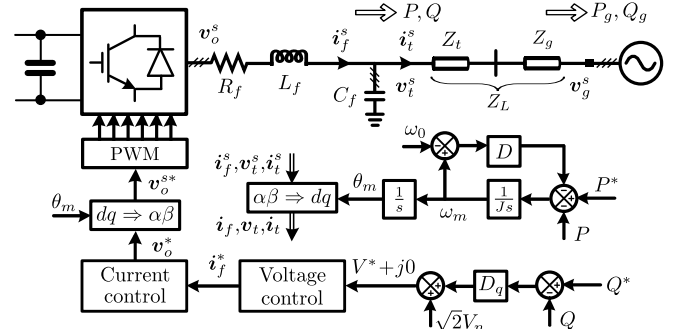


Fig. 1. Simplified diagram of a VSG-controlled converter.

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As depicted in Fig. 1, the active power transferred by the VSG is regulated by the swing equation as follows:

$$P^* - P = J \frac{d\omega_m}{dt} + D(\omega_m - \omega_0) \quad (1)$$

where  $J$  and  $D$  denote inertia and droop coefficients for active power regulation, respectively.  $\omega_m$  is the angular frequency of the VSG, and  $\omega_0$  represents the normal grid frequency.

From Fig. 1, VSG's instantaneous active and reactive power can be formulated as [13]

$$\begin{aligned} P &= 3/2 (v_{td}i_{td} + v_{tq}i_{tq}), \\ Q &= 3/2 (v_{tq}i_{td} - v_{td}i_{tq}) \end{aligned} \quad (2)$$

where  $v_{td}$ ,  $v_{tq}$ ,  $i_{td}$  and  $i_{tq}$  are capacitor voltage and grid-side current on the synchronous  $dq$  coordinates.

However, for modeling a VSG, (2) normally leads to a high-order model which includes all the dynamics of current and voltage loops [14], and this inevitably makes analysis and controller design complex. Without loss of generality, impacts of inner control loops can be neglected for simplicity, assuming that the bandwidth of these loops is much higher than that of power control loops [4], [8], [9], [15]. The VSG-controlled system is assumed to be balanced in most cases as well. Subsequently, the analysis can be conducted by using the phasor representation, where capacitor and grid voltage phasors are denoted by  $\dot{V}_t = V_t \angle \delta$  and  $\dot{V}_g = V_g \angle 0$ , respectively, and  $\delta$  is the power angle which can be expressed as

$$\delta = \int (\omega_m - \omega_g) dt \quad (3)$$

where  $\omega_g$  is the angular frequency of the grid voltage.

Assuming that the grid-side impedance is mainly inductive,  $Z_L$  can be then replaced by the reactance  $X_L$ . Note that, if the grid-side network is not mainly inductive, a virtual impedance loop can be used for adjusting the equivalent output impedance to be mainly inductive [16]. Active and reactive power of the VSG-controlled converter can be subsequently obtained from the phasor representation as [17]

$$\begin{aligned} P &= 3V_t V_g \sin(\delta) / X_L, \\ Q &= 3 [V_t^2 - V_t V_g \cos(\delta)] / X_L. \end{aligned} \quad (4)$$

Since  $\delta$  is normally small enough, approximation  $\sin \delta \approx \delta$  and  $\cos \delta \approx 1$  can be applied. In addition,  $V_t \approx V_g \approx V_n$ , where  $V_n$  denotes the rated phase voltage. It should be mentioned that, in the case of unbalanced conditions, (4) cannot be used to replace (2), a different modeling method should be applied.

With respect to the low-frequency power oscillation issue in the VSG-controlled converter, analyses are generally applied without considering the reactive power control since the oscillation is mainly introduced by the active power control [4], [8], [9], [18], i.e., using a large virtual inertia. In this manner, power loops are taken as fully decoupled, and only active power dynamics are included the analysis. By linearizing (1), (3), and (4), the small-signal model of the active power change  $\Delta P$  over the set point variation  $\Delta P^*$  is derived as [9], [15]

$$\Delta P = \frac{3V_n^2}{JX_L s^2 + DX_L s + 3V_n^2} \Delta P^*. \quad (5)$$

TABLE I  
VARIANTS OF THE SWING EQUATION

| Ref. | Active power control law                                                                                                                                                     | ADOF | $\frac{\Delta P}{\Delta P^*}$ |
|------|------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|------|-------------------------------|
| [6]  | $P^* - P = J \frac{d\omega_m}{dt} + D(\omega_m - \omega_0) + D_{pll}(\omega_m - \hat{\omega}_g)$                                                                             | 1    | $2^{nd}$                      |
| [7]  | $P^* - \frac{1+D}{1+T_f s} P = J \frac{d\omega_m}{dt} + D(\omega_m - \omega_0)$                                                                                              | 2    | $3^{th}$                      |
| [8]  | $P^* - \frac{1}{1+T_f s} P = J \frac{d\omega_m}{dt} + D(\omega_m - \omega_0) - P_d$<br>$P_d = k_{xp}(\omega_m - \omega_0) + \frac{k_{xp}}{1+T_f s} P + \frac{k_{xi}}{s} P_d$ | 4    | $4^{th}$                      |
| [9]  | $P^* - \left( \frac{k_{p1}s}{s+k_{p2}} + 1 \right) P = \left( J + \frac{k_{\omega 1}}{s+k_{\omega 2}} \right) \frac{d\omega_m}{dt} + D(\omega_m - \omega_0)$                 | 4    | $4^{th}$                      |
| [10] | $P^* - P = J \frac{d\omega_m}{dt} + D(\omega_m - \omega_0) + D_v \left( \frac{T_{\omega s}}{s+T_{\omega}} \right) \omega_m$                                                  | 2    | $3^{th}$                      |

### III. LOW-FREQUENCY POWER OSCILLATIONS DAMPING

From (5), it can be seen that power dynamics are determined once coefficient  $J$  is chosen for a desired inertia response and coefficient  $D$  is fixed according to the  $P$ - $f$  droop curve. Since there are no more degrees of freedom (DOF) for tuning the overall damping level, low-frequency power oscillations may appear in the VSG-controlled converter. For extra damping effects to attenuate oscillations, variants of the swing equation have been extensively proposed, as summarized in Table I.

It worth noting that only when the additional DOF (ADOF) introduced by extra damping terms is greater than or equal to the order of  $\Delta P / \Delta P^*$ , closed-loop poles can be freely adjusted as desired without changing coefficients  $J$  and  $D$ . Thus, only the variants developed in [8], [9] provide enough ADOF. Nevertheless, all the algorithms in Table I will introduce degradation of the inertial response [12]. For a better damping method, the requirements are twofold:

- 1) No degradation of the inertial response.
- 2) Provide enough ADOF for any desired damping level and power dynamic.

### IV. RFF-BASED DAMPING METHODS

To introduce enough damping effects on active power oscillations without affecting the inertial response, RFF control is utilized. As shown in Fig. 2, VSG's active power is regulated by both RFF terms and the classical swing equation.

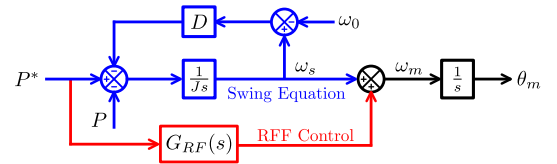


Fig. 2. Active power control with proposed RFF-based damping method.

From Fig. 2, VSG's new active power control law can be expressed as follows:

$$\begin{aligned} P^* - P &= J \frac{d\omega_s}{dt} + D(\omega_s - \omega_0), \\ \omega_m &= \omega_s + P^* G_{RFF}(s). \end{aligned} \quad (6)$$

where  $\omega_s$  is the swing-equation-regulated angular frequency.

In the steady state, for no impact on the static droop characteristic determined by the coefficient  $D$ , the RFF controller

$G_{RF}(s)$  should filter out most low frequency contents in the power reference, and one possible solution is as follows:

$$G_{RF1}(s) = k_{hp1} \frac{s}{s + k_{hp2}} \quad (7)$$

where  $k_{hp1}$  and  $k_{hp2}$  are two coefficients for adjusting the pass band of the RFF control loop. Combine (3), (4), (6) and (7), the following small-signal model can be derived:

$$\Delta P = \frac{3V_n^2 [Jk_{hp1}s^2 + (Dk_{hp1} + 1)s + k_{hp2}]}{(JX_L s^2 + DX_L s + 3V_n^2)(s + k_{hp2})} \Delta P^* \quad (8)$$

Here an additional real pole is introduced at  $-k_{hp2}$ . There are two extra zeros can be adjusted by both  $k_{hp1}$  and  $k_{hp2}$ . For adequate damping effects, two zeros need to be moved towards the poorly-damped poles. It is better to place the zeros along the natural frequency curve of those poles. In this way, the effect of zeros can mitigate the oscillations generated by the poorly-damped poles to the most extent. As shown in Fig. 3, a sweep of  $k_{hp1}$  is conducted for the best pole-zero placement.

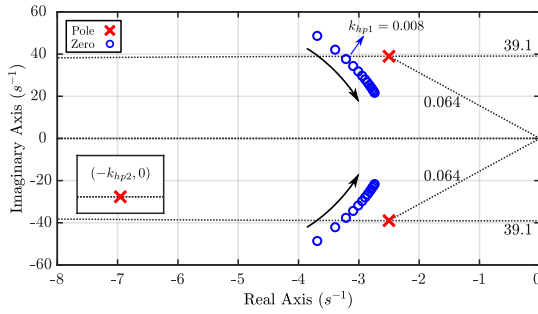


Fig. 3. Pole-zero map of  $\Delta P/\Delta P^*$ , where  $k_{hp1}$  varies from 0.006 to 0.03,  $k_{hp2} = 1000$ ,  $J = 70$ ,  $D = 350$ ,  $V_n = 380/\sqrt{3}$  V, and  $X_L = 1.35 \Omega$ .

Although poorly-damped oscillations can be eased by placing additional zeros, the ADOF is still limited for tuning the damping effect. Therefore, for adjusting power dynamics to any desired level, controller  $G_{RF}(s)$  can be directly derived by assuming  $\Delta P/\Delta P^*$  is a second order plant as

$$\Delta P = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2} \Delta P^* \quad (9)$$

where  $\zeta$  is the desired damping factor, and  $\omega_n$  is the desired natural frequency. From (3), (4), (6) and (9), the RFF controller  $G_{RF}(s)$  can be directly derived as follows:

$$G_{RF2}(s) = \frac{m_2 s^2 + m_1 s}{3V_n^2 (Js^3 + n_2 s^2 + n_1 s + D\omega_n^2)} \quad (10)$$

where the coefficients in (10) are

$$\begin{aligned} m_2 &= J\omega_n^2 X_L - 3V_n^2, & m_1 &= D\omega_n^2 X_L - 6V_n^2 \zeta \omega_n, \\ n_2 &= D + 2J\zeta \omega_n, & n_1 &= J\omega_n^2 + 2D\zeta \omega_n. \end{aligned} \quad (11)$$

From (11), it can be seen that all the controller coefficients are calculated directly according to parameters of the VSG-controlled system, damping factor  $\zeta$  and natural frequency  $\omega_n$ , which significantly benefits the controller automatic tuning and real-time adaptation. With respect to the selection of  $\zeta$  and  $\omega_n$ , the classical criteria which is normally used for tuning second-order plants can be adopted. For well-damped active power

responses, the damping factor  $\zeta$  can be set to a value greater than 0.7. Afterwards, the natural frequency  $\omega_n$  can be selected according to the settling time requirement, i.e.,  $\omega_n = T_{set}\zeta/4$ , where  $T_{set}$  is the settling time [19].

It should be mentioned that the aforementioned theoretical analysis is conducted under the assumption that active and reactive power loops of a VSG are fully decoupled; however, it cannot be ensured in some cases [15]. Therefore, it is necessary to assess the coupling effect of VSG's reactive power loop on the proposed scheme. Taking the coupling effect into account, the small-signal model of the active power loop can be written as

$$\Delta P = \frac{3V_n^2 [1 + (Js + D)G_{RF}(s)]T_c}{JX_L s^2 + DX_L s + 3V_n^2 T_c} \Delta P^* \quad (12)$$

where  $T_c$  represents the coupling effect which is introduced by the reactive power loop. The derivation of (12) and  $T_c$  are given in the Appendix.

Using (7), (10) and (12), the impact of the power-loop coupling can be assessed by the pole-zero map of  $\Delta P/\Delta P^*$ . To have a convincing evaluation of the coupling effect, the static power angle  $\delta_n$  in (A.3) is firstly set to a large value, for instance, 1 rad. The droop gain  $D_q$  in (A.2) varies from 0 pu to 10 pu. As shown in Fig. 4, either RFF controller  $G_{RF1}(s)$  or  $G_{RF2}(s)$  is applied, the coupling effect on the effectiveness of improving the overall damping level can be neglected since the desired pole-zero placement is barely changed by  $D_q$ .

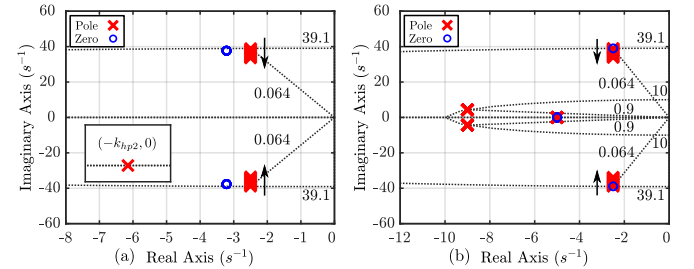


Fig. 4. Pole-zero map of  $\Delta P/\Delta P^*$  when  $\delta_n = 1$  rad, and  $D_q$  changes from 0 pu to 10 pu.  $J = 70$ ,  $D = 350$ ,  $V_n = 380/\sqrt{3}$  V, and  $X_L = 1.35 \Omega$ . (a) Controller  $G_{RF1}(s)$  is applied with  $k_{hp1}=0.008$  and  $k_{hp2}=1000$ ; (b) Controller  $G_{RF2}(s)$  is applied with  $\zeta=0.9$  and  $\omega_n=10$  rad/s.

## V. SIMULATION AND EXPERIMENTAL VALIDATIONS

To validate the effectiveness of proposed RFF-based damping method, EMT simulations based on the averaged value of converters have been conducted in Digsilent/PowerFactory, and a laboratory implementation as Fig. 5 is performed as well. In the studies, two cases shown in Fig. 6 have been included, and the active power is obtained via the instantaneous calculation as (2). All the parameters used in the studies are listed in Table II.

### A. Damping of low-frequency oscillations - Simulation

In the simulation studies, the grid-tied mode is tested firstly. A step change from 0 to 0.6 pu is applied in the active power set point. As shown in Fig. 7(a), without any additional damping terms, output oscillates at a low frequency (around

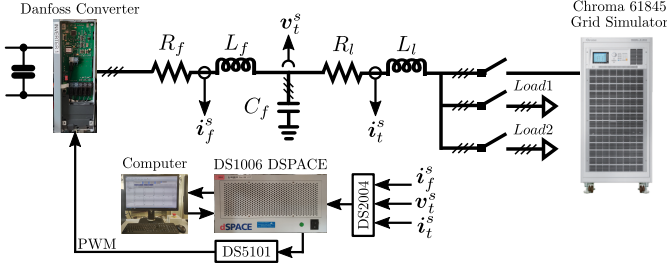


Fig. 5. Experiment setup.

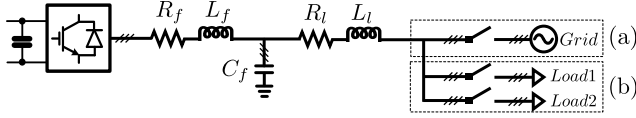


Fig. 6. Circuit diagrams of two cases. (a) Damping of low-frequency power oscillations in grid-tied mode; (b) Inertial response in islanded mode.

6 Hz). However, when  $G_{RF1}(s)$  is applied with  $k_{hp1}=0.008$  and  $k_{hp2}=1000$ , damping effects are significantly improved. Especially, with  $G_{RF2}(s)$  applied with  $\zeta = 0.9$  and  $\omega_n = 10$  rad/s, poorly-damped oscillations are attenuated to zero.

### B. The inertial response - Simulation

To check the inertial response, the approach used in [4] and the RoCoF calculation method specified in [20] are adopted. The VSG initially supplies 0.6 kW local load, and a step change of 0.6 kW is suddenly applied in the local load to simulate the load variation. The resulting angular frequencies are shown in Fig. 7(b). It can be seen that angular frequencies change in the same way, with or without the proposed RFF-based damping method, and the values of RoCoF are approximately the same. The results indicate that the inertial response of VSG is well preserved.

### C. Damping of low-frequency oscillations - Experiment

A step change of 0.6 kW is applied in the active power set point of grid-tied VSG to demonstrate power oscillations. As

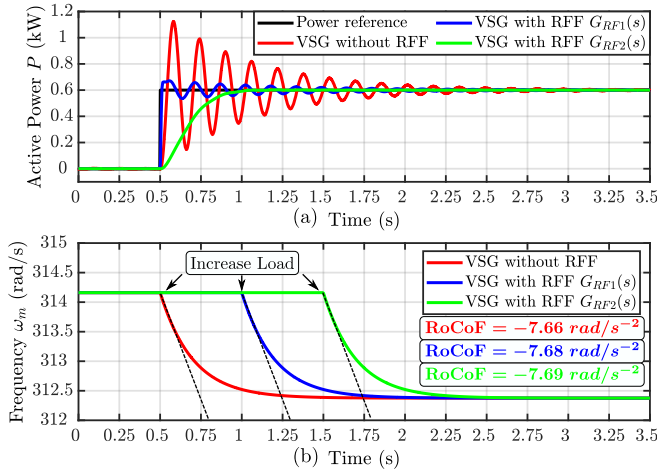


Fig. 7. Simulation results (a) VSG output active power in grid-tied mode; (b) VSG angular frequency in islanded mode.

TABLE II  
PARAMETERS OF THE VSG-CONTROLLED CONVERTER

| Symbol     | Values simulation | Values experiment | Symbol     | Values simulation | Values experiment |
|------------|-------------------|-------------------|------------|-------------------|-------------------|
| $S_{base}$ | 2.2 kVA           | 2.2 kVA           | $V_n$      | $380/\sqrt{3}$ V  | $380/\sqrt{3}$ V  |
| $\omega_0$ | 314 rad/s         | 314 rad/s         | $V_{dc}$   | 650 V             | 650 V             |
| $R_f$      | 0 $\Omega$        | 0.1 $\Omega$      | $L_f$      | 3.6 mH            | 3.6 mH            |
| $C_f$      | 9 $\mu$ F         | 9 $\mu$ F         | $R_l$      | 0 $\Omega$        | 0.5 $\Omega$      |
| $L_l$      | 4.3 mH            | 4.3 mH            | $f_{sw}$   | N/A               | 10 kHz            |
| $J$        | 70 (10 pu)        | 70 (10 pu)        | $D$        | 350 (50 pu)       | 350 (50 pu)       |
| $D_q$      | 0.014 (0.1 pu)    | 0.014 (0.1 pu)    |            |                   |                   |
| $k_{hp1}$  | 0.008             | 0.03              | $k_{hp2}$  | 1000              | 1000              |
| $\zeta$    | 0.9               | 0.9               | $\omega_n$ | 10 rad/s          | 10 rad/s          |

shown in Fig 8(a), after the set point change, oscillations at approximately 4 Hz appear in the output power. Because of inevitable mismatches in the setup, for example, impedance variations, sampling errors and nonlinear properties of the power converter, the oscillation frequency is not exactly the same as that in the simulation. The difference is around 2 Hz. However, it does not affect the experimental assessment.

As shown in Fig. 8(b), when the  $G_{RF1}(s)$  is applied with  $k_{hp1}=0.03$  and  $k_{hp2}=1000$ , low-frequency power oscillations can be significantly attenuated as the simulation result shown in Fig. 7. The difference between simulation ( $k_{hp1} = 0.008$ ) and experimental ( $k_{hp1} = 0.03$ ) settings reveals that controller  $G_{RF1}(s)$  is sensitive to the parameter variations. Thus, for appropriate damping effects, an accurate placement of zeros introduced by  $G_{RF1}(s)$  is essential. In this manner, an optimization method may be required in real applications when  $G_{RF1}(s)$  is applied to VSGs. On the other hand, experimental results presented in Fig. 8(c) show that controller  $G_{RF2}(s)$  is still effective under the same condition. The coefficients  $\zeta$  and  $\omega_n$  are set to 0.9 and 10 rad/s in both simulation and experiment. As depicted in Fig. 8(c), smooth power dynamics without any overshoot and oscillation can be attained, which is the same as the result obtained in the simulation studies.

### D. The inertial response - Experiment

The VSG operates in islanded mode, and one resistor bank of 230  $\Omega$  is used as the local load. Another 230  $\Omega$  resistor bank is switched on to emulate a sudden load variation, just as the scenario applied in [4]. As shown in Fig. 9, with or without using the proposed RFF-based damping method, angular frequencies of VSG change in the same profile. Especially, the values of RoCoF are almost the same, which shows that the inertial response of VSG is not changed.

### E. Response to the grid frequency variation - Experiment

For a further assessment of the proposed control scheme under grid frequency variations, a step frequency change of 0.2 Hz is applied in the experiment. VSGs are connected to the grid, and the grid frequency drops suddenly from the nominal value 50 Hz to 49.8 Hz. As shown in Fig. 10, responses of grid-tied VSGs with or without feedforward controllers are approximately the same. Specifically, at the moment of frequency step change, power angle  $\delta$  in (4) starts to increase,



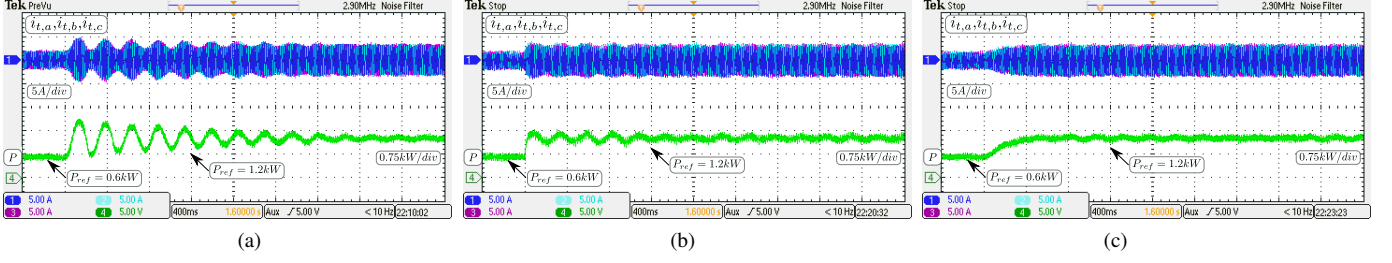


Fig. 8. Waveforms of the grid-tied VSG. (a) Without proposed damping methods; (b) Controller  $G_{RF1}(s)$  is applied; (c) Controller  $G_{RF2}(s)$  is applied.

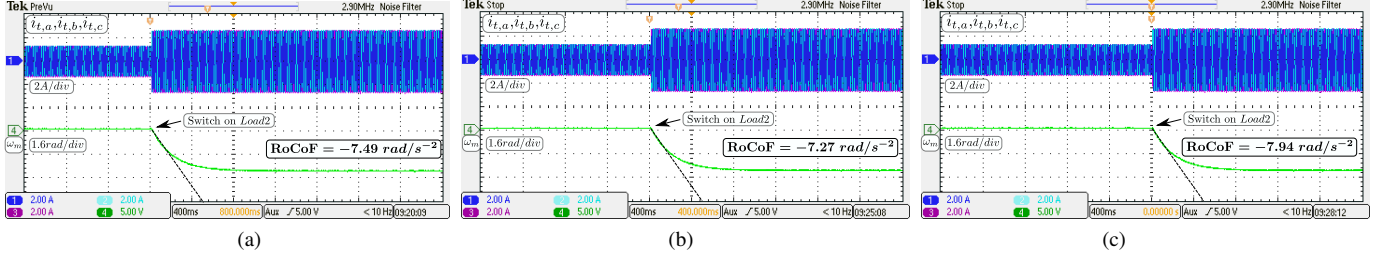


Fig. 9. Waveforms of the islanded VSG. (a) Without proposed damping methods; (b) Controller  $G_{RF1}(s)$  is applied; (c) Controller  $G_{RF2}(s)$  is applied.

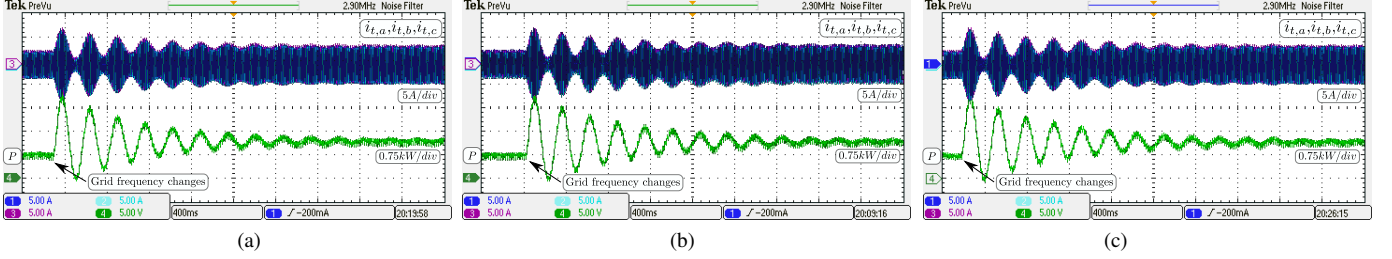


Fig. 10. Waveforms of the grid-tied VSG under grid frequency variations. (a) Without proposed damping methods; (b) Controller  $G_{RF1}(s)$  is applied; (c) Controller  $G_{RF2}(s)$  is applied.

and the output power rise rapidly to inject inertial power for supporting the grid frequency. Afterwards, because original inertial responses of tested VSGs are preserved without any degradation, the output power swings until a new operation point is reached.

## VI. CONCLUSION

This paper proposes a simple but effective method to actively attenuate the poorly-damped power oscillations in VSG-controlled converters. Different from the existing algorithms using the feedback control, the proposed method introduces a feed-forward path to improve the damping level. By applying the proposed method, threefold benefits have been attained: i) Power oscillations can be effectively attenuated without affecting the original inertial response; ii) The damping effect and power dynamics can be flexibly adjusted to any specific level; iii) Controller automatic tuning and real-time adaptation are enabled, provided that system parameters are available.

It is worthy to note that, under sudden grid frequency variations, the rapid active power change which comes from the original inertial response is favorable for maintaining the grid frequency stability; nevertheless, the resulting power swings are not necessary. To handle this issue in grid-tied VSGs, additional control strategies could be developed and implemented with the proposed control scheme.

## APPENDIX

From (4), active and reactive power variations, which include the coupling items, can be derived as [15]

$$\begin{aligned}\Delta P &= 3V_n^2 \Delta\delta / X_L + 3V_n (\Delta V_t + \Delta V_g) \delta_n / X_L, \\ \Delta Q &= 3V_n (\Delta V_t - \Delta V_g) / X_L + 3V_n^2 \delta_n \Delta\delta / X_L\end{aligned}\quad (\text{A.1})$$

where  $\delta_n$  denotes the static power angle.

In (A.1),  $\Delta V_g \approx 0$ , and  $\Delta V_t$  can be formulated as

$$\Delta V_t = D_q (\Delta Q^* - \Delta Q) \quad (\text{A.2})$$

where  $D_q$  denotes the droop gain.

Taking both (A.2) and (A.1) into account, the active power variation, which contains the reactive power loop gain, can be rewritten as

$$\Delta P = \frac{3V_n^2}{X_L} \underbrace{\left(1 - \frac{T_q}{1 + T_q} \delta_n^2\right)}_{T_c} \Delta\delta \quad (\text{A.3})$$

where  $T_q$  is the reactive power loop gain which is equal to  $3V_n D_q / X_L$ . From (3), (6) and (A.3), the active power change  $\Delta P$  over the set point variation  $\Delta P^*$  can be derived as (12).

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